## CHAPTER 7

## Sinusoids and Phasors

Recall that, for capacitors and inductors, the branch variables (current values and voltage values) are related by differential equations. Normally, to analyze a circuit containing capacitor and/or inductor, we need to solve some differential equations. The analysis can be greatly simplifies when the circuit is driven (or excited) by a source (or sources) that is sinusoidal. Such assumption will be the main focus of this chapter.

### 7.1. Prelude to Second-Order Circuits

The next example demonstrates the complication normally involved when analyzing a circuit containing capacitor and inductor. This example and the analysis presented is not the main focus of this chapter.

Example 7.1.1. The switch in the figure below has been open for a long time. It is closed at $t=0$.

$$
v=L \frac{d i}{d t}
$$


(a) Find $v(0)$ and $\frac{d v}{d t}(0)$
(b) Find $v(t)$ for $t>0$. $\}$ The source is $v_{s}(t) \equiv 12 \mathrm{~V}$
(c) Find $v(\infty)$ and $\frac{d v}{d t}(\infty)$.
(d) Find $v(t)$ for $t>0$ when the source is $v_{s}(t)= \begin{cases}12, & t<0, \\ 12 \cos (t), & t \geq 0 .\end{cases}$


At $t=0, i(0)=0 A(n o$ jump in $i$ through $L$ )
$V(0)=12 \mathrm{~V}$ (no jump in $v$ across $C$ )
Due to the we only have to focus on the loop on the RHS. o A source $92 \geq$,
which is equivalent, to an open circuit,
 this part of ' the circuit has no current and hence it is just a hanging branch.
(b) For $t>0$,


$$
\begin{array}{rlr}
K C L @ A: \quad-i+\frac{v}{R_{2}}+i_{c} & =0 & i_{c}=c \frac{d v}{d t} \\
\Rightarrow i & =\frac{v}{R_{2}}+i_{c}=\frac{v}{R_{2}}+c v^{\prime}
\end{array}
$$

Also, $\quad i_{c}=c \frac{d v}{d t}$.

$$
\text { Therefore, } \quad \frac{d v}{d t}=\frac{1}{c} i_{c}=\frac{-6}{1 / 2}=-12 \mathrm{~V} / \mathrm{s} \text {. }
$$



$$
\frac{d v}{d t}(0)=-12 \mathrm{~V} / \mathrm{s}
$$

KVL around the outer-loop:

$$
\begin{array}{r}
+v_{s}-i R_{1}-v_{L}-v=0 \\
v_{L}=L \frac{d i}{d t} \\
v_{1}-i R-v=0
\end{array}
$$

$$
\Rightarrow+v_{s}-i R_{1}-L i i^{\prime}-v=0
$$

$$
v_{s}-\left(\frac{v}{R_{2}}+c v^{\prime}\right) R_{1}-L\left(\frac{v}{R_{2}}+c v^{\prime}\right)^{\prime}-v=0
$$

$$
v^{\prime \prime}+\left(\frac{1}{C R_{2}}+\frac{R_{1}}{L}\right) v^{\prime}+\frac{1}{L C}\left(1+\frac{R_{1}}{R_{2}}\right) v=\frac{v_{s}}{L C}
$$

$$
\mathrm{v}=\mathrm{dsolve}\left(\mathrm{D} 2 \mathrm{v}+5 * \mathrm{Dv}+6 * \mathrm{v}=24 \mathrm{I}^{\prime}, \mathrm{v}(0)=12^{\prime}, \mathrm{Dv}(0)=-12 \prime\right)
$$

gives $v(t)=4+12 e^{-2 t}-4 e^{-3 t}$. Similarly,
(c) $v(\infty)=4, \frac{d v}{d t}(\infty)=0$
(d) $\mathrm{v}=$ dsolve('D2v + 5*Dv + 6*v $=2 * 12 * \cos (\mathrm{t}) \mathrm{I}, \mathrm{v}(0)=12$ ', $\mathrm{Dv}(0)=-12$ ','t')
gives $v(t)=\frac{72}{5} e^{-2 t}-\frac{24}{5} e^{-3 t}+\frac{12}{5} \cos (t)+\frac{12}{5} \sin (t) ., \quad t>0$



$$
\Leftrightarrow(j \omega)^{2} \vec{V}+5 j \omega \stackrel{L}{\rightharpoonup} \vec{V}+6 \vec{V}=24 L 0^{0}
$$

$$
\begin{aligned}
&-\vec{v}+5 j \vec{v}+6 \vec{v}=24 \\
& \stackrel{\rightharpoonup}{v}=\frac{24}{5+5 j}=\frac{12 \sqrt{2}}{5} L-45^{\circ} \\
& \Leftrightarrow v(t)=\frac{12 \sqrt{2}}{5} \cos \left(t-45^{\circ}\right)
\end{aligned}
$$

### 7.2. Sinusoids

Definition 7.2.1. Some terminology:
(a) A sinusoid is a signal (, e.g. voltage or current) that has the form of the sine or cosine function. (with o average)

- Turn out that you can express them all under the same notation using only cosine (or only sine) function.
- We will use cosine.
(b) A sinusoidal current is referred to as alternating current (AC).
(c) We use the term AC source for any device that supplies a sinusoidally varying voltage (potential difference) or current.
(d) Circuits driven by sinusoidal current or voltage sources are called AC circuits.
7.2.2. Consider the sinusoid 21 signal (in cosine form) current $i(t)$ where

$$
\begin{aligned}
& \text { angular freq. freq. } \\
& \text { dah signal (in cosine form) } \\
& 2 t+\phi)=X_{m} \cos (2 \pi f t+\phi) \text {, }
\end{aligned}
$$

$X_{m}$ : the amplitude of the sinusoid,
$2 \pi f=\omega$ : the angular frequency in radians/s (or rad /s),
$\phi$ : the phase.

- First, we consider the case when $\phi=0$ :
- When $\phi \neq 0$, we shift the graph of $X_{90^{\circ}} X_{m} \cos (\omega t)$ to the left "by $\phi$ ".

7.2.3. The period (the time of one complete cycle) of the sinusoid is

$$
T=\frac{2 \pi}{\omega}
$$

The unit of the period is in second if the angular frequency unit is in radian per second.

The frequency $f$ (the number of cycles per second or hertz ( Hz ) ) is the reciprocal of this quantity, i.e.,

$$
f=\frac{1}{T}
$$

7.2.4. Standard form for sinusoid: In this class, when you are asked to find the sinusoid representation of a signal, make sure that your answer is in the form

$$
x(t)=X_{m} \cos (\omega t+\phi)=X_{m} \cos (2 \pi f t+\phi)
$$

where $X_{m}$ is nonnegative and $\phi$ is between $-180^{\circ}$ and $+180^{\circ}$

### 7.2.5. Conversions to standard form

- When the signal is given in the sine form, it can be converted into its cosine form via the identity $=\cos \left(x+\left(-90^{\circ}\right)\right) \equiv \cos$ "shifted to

$$
\sin (x)=\cos \left(x-90^{\circ}\right)
$$


the left by $-90^{\circ}=$ ミcos"shifted to the right by $90^{\circ}$

In particular,

$$
X_{m} \sin (\omega t+\phi)=X_{m} \cos \left(\omega t+\phi-90^{\circ}\right)
$$

- $X_{m}$ is always non-negative. We can avoid having the negative sign by the following conversion:

$$
\begin{aligned}
-\cos (x)=\cos \left(x \pm 180^{\circ}\right) . \\
\operatorname{Eos}\left(x+110^{\circ}\right)
\end{aligned}
$$

In particular,

$$
-A \cos (\omega t+\phi)=A \cos \left(2 \pi f t+\phi \pm 180^{\circ}\right)
$$

Note that usually you do not have the choice between $+180^{\circ}$ or $-180^{\circ}$. The one that you need to use is the one that makes $\phi \pm 180^{\circ}$ falls somewhere between $-180^{\circ}$ and $+180^{\circ}$.
7.2.6. For any ${ }^{11}$ linear AC circuit, the "steady-state" voltage and current are sinusoidal with the same frequency as the driving source(s).

- Although all the voltage and current are sinusoidal, their amplitudes and phases can be different.
- These can be found by the technique discussed in this chapter.


### 7.3. Phasors

Sinusoids are easily expressed in terms of phasors, which are more convenient to work with than sine and cosine functions. The tradeoff is that phasors are complex-valued.

$$
i=\sqrt{-1}
$$

7.3.1. The idea of phasor representatiifn is based on Euler's identity:

$$
e^{j \phi}=\cos \phi+j \sin \phi, \quad \varnothing \text { is any real number }
$$

From the identity, we may regard $\cos \phi$ and $\sin \phi$ as the real and imaginary parts of $e^{j \phi}$ :

$$
\cos \phi=\operatorname{Re}\left\{e^{j \phi}\right\}, \quad \sin \phi=\operatorname{Im}\left\{e^{j \phi}\right\},
$$

where Re and Im stand for "the real part of" and "the imaginary part of" $e^{j \phi}$.

Definition 7.3.2. A phasor is a complex number that represents the amplitude and phase of a sinusoid. Given a sinusoid $x(t)=X_{m} \cos (\omega t+\phi)$, then


The complex number $\mathbf{X}$ is called the phasor representation of the sinusoid $v(t)$. Notice that a phasor captures information about amplitude and phase of the corresponding sinusoid.

$$
\text { Ex. } x(t)=\sqrt{2} \cos \left(10 \pi t+45^{\circ}\right) \Leftrightarrow \vec{x}=\sqrt{2}<45^{\circ}
$$

[^0]7.3.3. Whenever a sinusoid is expressed as a phasor, the term $e^{j \omega t}$ is implicit. It is therefore important, when dealing with phasors, to keep in mind the frequency $f$ (or the angular frequency $\omega$ ) of the phasor.
7.3.4. Given a phasor $\mathbf{X}$, to obtain the time-domain sinusoid caresponging to a given phasor, there are two important routes.
(a) Simply write down the cosine function with the same magnitude as the phasor and the argument as $\omega t$ plus the phase of the phasor.
(b) Multiply the phasor by the time factor $e^{j \omega t}$ and take the real part.
7.3.5. Any complex number $z$ (including any phasor) can be equivalently represented in three forms.
(a) Rectangular form: $z=x+j y$.
(b) Polar form: $z=r \angle \phi$.
(c) Exponential form: $z=r e^{j \phi}$
where the relations between them are
\[

$$
\begin{gathered}
r=\sqrt{x^{2}+y^{2}}, \quad \phi=\tan ^{-1} \frac{y}{x} \pm 180^{\circ} . \\
x=r \cos \phi, \quad y=r \sin \phi .
\end{gathered}
$$
\]

Note that for $\phi$, the choice of using $+180^{\circ}$ or $-180^{\circ}$ in the formula is determined by the actual quadrant in which the complex number lies.


As a complex quantity, a phasor may be expressed in rectangular form, polar form, or exponential form. In this class, we focus on polar form.
7.3.6. Summary: By suppressing the time factor, we transform the sinusoid from the time domain to the phasor domain. This transformation is summarized as follows: same convention

$$
x(t)=X_{m} \cos (\omega t+\phi) \stackrel{\vdots}{\Leftrightarrow} \mathbf{X}=X_{m} \angle \phi . \quad x_{m} \geqslant 0
$$

Time domain representation $\Leftrightarrow$ Phasor domain representation $180^{\circ}<\phi \leqslant 180^{\circ}$ Real-valued $A$ complex Function of time: A complex number
(time-dependent)

Give instantaneous value : Frequency domain of the signal at tine $t$. 1

DEFINITION 7.3.7. Standard form for phasor: In this class, when you are asked to find the phasor representation of a signal, make sure that your answer is a complex number in polar form, i.e. $r \angle \phi$ where $r$ is nonnegative and $\phi$ is between $-180^{\circ}$ and $+180^{\circ}$.

Example 7.3.8. Transform these sinusoids to phasors:
(a) $i=6 \cos \left(50 t-40^{\circ}\right) \mathrm{A} \Leftrightarrow \vec{I}=6 \angle-40^{\circ}$
(b) $v=-4 \sin \left(30 t+50^{\circ}\right) \mathrm{V}=-4 \cos \left(30 t+50^{\circ}-90^{\circ}\right)$

$$
\begin{aligned}
& =-4 \cos \left(30 t-40^{\circ}\right)=4 \cos \left(30 t-40^{\circ} \pm 180^{\circ}\right) \\
& =4 \cos \left(30 t+140^{\circ}\right) \Leftrightarrow \vec{V}=4<140^{\circ}
\end{aligned}
$$

EXAMPLE 7.3.9. Find the sinusoids represented by these phasors:
(a) $\mathbf{I}=-3+j 4 A=5 \angle 127^{\circ} \Leftrightarrow i(t)=5 \cos \left(\omega t+127^{\circ}\right)$ call
(b) $\mathbf{V}=(j) 3 e^{-j 20^{\circ}} \mathrm{V}^{\circ}=\left(1 \angle 90^{\circ}\right)\left(8 \angle-20^{\circ}\right)=8 \angle 70^{\circ} \Leftrightarrow v(t)=$ $e^{j 90^{\circ}} \times 8 \times e^{-j 10^{\circ}} \quad 8 \cos (\omega t$

$$
\left.\left(r_{1} L \theta_{1}\right)\left(r_{2}<\theta_{2}\right)=r_{1} r_{2}<\left(\theta_{1}+\theta_{2}\right)+70^{\circ}\right)
$$

7.3.10. The differences between $x(t)$ and $\mathbf{X}$ should be emphasized:
(a) $x(t)$ is the instantaneous or time-domain representation, while $\mathbf{X}$ is the frequency or phasor-domain representation.
(b) $x(t)$ is time dependent, while $\mathbf{X}$ is not.
(c) $x(t)$ is always real with no complex term, while $\mathbf{X}$ is generally complex.
7.3.11. Adding sinusoids of the same frequency is equivalent to adding their corresponding phasors. To see this,

$$
\begin{aligned}
A_{1} \cos \left(\omega t+\phi_{1}\right)+A_{2} \cos \left(\omega t+\phi_{2}\right) & =\operatorname{Re}\left\{\mathbf{A}_{1} e^{j \omega t}\right\}+\operatorname{Re}\left\{\mathbf{A}_{2} e^{j \omega t}\right\} \\
& =\operatorname{Re}\left\{\left(\mathbf{A}_{1}+\mathbf{A}_{2}\right) e^{j \omega t}\right\}
\end{aligned}
$$

Because $\mathbf{A}_{1}+\mathbf{A}_{2}$ is just another complex number, we can conclude a (surprising) fact: adding two sinusoids of the same frequency gives another sinusoids. $\cos (t-90)$
Ex. $\cos (t)+\widetilde{\sin (t)}$
$\Leftrightarrow 1 \angle 0^{\circ}+1 \angle-90^{\circ}=\sqrt{2} \angle-45^{\circ}$
$\Leftrightarrow \sqrt{2} \cos \left(1 t-45^{\circ}\right)$


$$
\omega=2
$$

98
7. SINUSOIDS AND PHASORS
$\left.=4 \cos \left(2 t+0^{\circ}\right)+3 \cos (3) t-90^{\circ}\right)$
Example 7.3.12. $x(t)=4 \cos (2 t)+3 \sin (2 t)$
Direct calculation:
$4 \cos \left(2 t+0^{\circ}\right)+3 \cos \left(2 t-90^{5}\right)$ 告) $+\cos (2 t)$
$=\operatorname{Re}\left\{\left(4 \angle 0^{\circ}\right) e^{j \angle t}\right\}$
$=\operatorname{Re}\left\{\left(4 \angle 0^{\circ}\right)+(3 L-90) e^{j 2 t}\right\}=V^{j 2} V_{-1} . V_{-i s} . V_{i s}=5 L-36.87^{\circ}$
$=\operatorname{Re}\left\{\left(5 \angle-36.87^{\circ}\right) e^{j 2 t}\right\}$
$\Leftrightarrow 5 \cos \left(2 t-36.87^{\circ}\right)$
7.3.13. Properties involving differentiation and integration:
$=5 \cos (2 t$ (a) Differentiating a sinusoid is equivalent to multiplying its care-

$$
\frac{d x(t)}{d t} \Leftrightarrow j \omega \mathbf{X} .
$$

To see this, suppose $x(t)=X_{m} \cos (\omega t+\phi)$. Then,

$$
\begin{aligned}
\frac{d x}{d t}(t) & =-\omega X_{m} \sin (\omega t+\phi)=\omega X_{m} \cos \left(\omega t+\phi-90^{\circ}+180^{\circ}\right) \\
& =\operatorname{Re}\left\{\omega X_{m} e^{j \phi} e^{j 90^{\circ}} \cdot e^{j \omega t}\right\}=\operatorname{Re}\left\{j \omega \mathbf{X} e^{j \omega t}\right\}
\end{aligned}
$$

Alternatively, express $v(t)$ as

$$
x(t)=\operatorname{Re}\left\{X_{m} e^{j(\omega t+\phi)}\right\} .
$$

Then,

$$
\frac{d}{d t} x(t)=\operatorname{Re}\left\{X_{m} j \omega e^{j(\omega t+\phi)}\right\}
$$

(b) Integrating a sinusoid is equivalent to dividing its corresponding phasor by $j \omega$. In other words,

$$
\int x(t) d t \Leftrightarrow \frac{\mathbf{X}}{j \omega} .
$$

Example 7.3.14. Find the voltage $v(t)$ in a circuit described by the intergrodifferential equation

$$
\left.2 \frac{d v}{d t}+5 v+10 \int v d t=50 \cos (5)-30^{\circ}\right) \quad \rightarrow^{w=5} \quad \frac{2}{j} \times \frac{j}{j}=\frac{2 j}{-1}=-2 j
$$

using the phasor approach.
$\Leftrightarrow$
(phasor domain)

$$
\begin{aligned}
& 2 j \omega \vec{V}+5 \vec{V}+10 \underset{j \omega}{\vec{V}}=50 \angle-30^{\circ} \left\lvert\, \vec{V}=\frac{50 L-30^{\circ}}{10 j+5-2 j}\right. \\
& =5.3 \mathrm{~L}-88^{\circ}
\end{aligned}
$$

$$
\vec{v}\left(10 j+5+\frac{2}{j}\right)=50 L-30^{\circ} \left\lvert\, \begin{aligned}
& \Leftrightarrow v(t)=5.3 \cos (5 t \\
& \stackrel{y}{\text { steady -state }} \quad
\end{aligned}\right.
$$

### 7.4. Phasor relationships for circuit elements - $C$

7.4.1. Resistor $R$ : If the current through a resistor $R$ is

$$
i(t)=I_{m} \cos (\omega t+\phi) \Leftrightarrow \mathbf{I}=I_{m} \angle \phi,
$$

the voltage across it is given by ohm's law


The phasor of the voltage is

$$
\mathbf{V}=R I_{m} \angle \phi .
$$

Hence,

$$
\mathbf{V}=\mathbf{I} R .
$$

We note that voltage and current are in phase and that the voltage-current relation for the resistor in the phasor domain continues to be Ohms law, as in the time domain.

7.4.2. Capacitor $C$ : If the voltage across a capacitor $C$ is

$$
v(t)=V_{m} \cos (\omega t+\phi) \Leftrightarrow \mathbf{V}=V_{m} \angle \phi,
$$

the current through it is given by

$$
i(t)=C \frac{d v(t)}{d t} \Leftrightarrow \mathbf{I}=j \omega C \mathbf{V}=\omega C V_{m} \angle\left(\phi+90^{\circ}\right) .
$$

The voltage and current are $90^{\circ}$ out of phase. Specifically, the current leads the voltage by $90^{\circ}$.


- Mnemonic: CIVIL

In a Capacitive (C) circuit, I leads V. In an inductive (L) circuit, V leads I.
7.4.3. Inductor $L$ : If the current through an inductor $L$ is

$$
i(t)=I_{m} \cos (\omega t+\phi) \Leftrightarrow \mathbf{I}=I_{m} \angle \phi,
$$

the voltage across it is given by

$$
v(t)=L \frac{d i(t)}{d t} \Leftrightarrow \mathbf{V}=j \omega L \mathbf{I}=\omega L I_{m} \angle\left(\phi+90^{\circ}\right) .
$$

$$
v=L \frac{d i}{d t}
$$


$\mathbf{V}=j \omega L \mathbf{I}$
The voltage and current are $90^{\circ}$ out of phase. Specifically, the current lags the voltage by $90^{\circ}$.



| Element | Time domain | Frequency domain | $z$ |
| :---: | :---: | :---: | :---: |
| $R$ | $v=R i$ | $\mathbf{V}=R \mathbf{I} 1 \angle 90^{\circ}$ | $R$ |
| $L$ | $v=L \frac{d i}{d t}$ | $\mathbf{V}=(j \omega L \mathbf{I}$ | $j \omega L$ |
| C | $i=C \frac{d v}{d t}$ | $\mathbf{V}=\frac{\mathbf{I}}{j \omega C}$ | $\frac{1}{j \omega c}=-\frac{1}{\omega c} j$ |

### 7.5. Impedance and Admittance

In the previous part, we obtained the voltage current relations for the three passive elements as

$$
\mathbf{V}=\mathbf{I} R, \quad \mathbf{V}=j \omega L \mathbf{I}, \quad \mathbf{I}=j \omega C \mathbf{V} .
$$

These equations may be written in terms of the ratio of the phasor voltage to the phasor of current as

$$
\frac{\mathbf{V}}{\mathbf{I}}=R, \quad \frac{\mathbf{V}}{\mathbf{I}}=j \omega L, \quad \frac{\mathbf{V}}{\mathbf{I}}=\frac{1}{j \omega C} \cdot \times \frac{j}{j}=j\left(-\frac{1}{\omega c}\right)
$$

From these equations, we obtain Ohm's law in phasor form for any type of element as

$$
\mathbf{Z}=\frac{\mathbf{V}}{\mathbf{I}} \quad \text { or } \quad \mathbf{V}=\mathbf{I Z}
$$

Definition 7.5.1. The impedance $\mathbf{Z}$ of a circuit is the ratio of the phasor voltage $\mathbf{V}$ to the phasor current $\mathbf{I}$, measured in ohms $(\Omega)$.

As a complex quantity, the impedance may be expressed in rectangular form as
with

$R=\operatorname{Re}\{\mathbf{Z}\}$ is called the resistance and $X=\operatorname{Im}\{\mathbf{Z}\}$ is called the reactance.

The reactance $X$ may be positive or negative. We say that the impedance is inductive when $X$ is positive or capacitive when $X$ is negative.

Definition 7.5.2. The admittance $(\mathbf{Y})$ is the reciprocal of impedance, measured in Siemens (S). The admittance of an element(or a circuit) is the ratio of the phasor current through it to phasor voltage across it, or

$$
\mathrm{Y}=\frac{\mathbf{1}}{\mathrm{Z}}=\frac{\mathbf{I}}{\mathrm{V}}
$$

$$
\vec{V}=\vec{I} \vec{z}
$$

7.5. IMPEDANCE AND ADMITTANCE
7.5.3. Kirchhoff's laws (KCL and KVL) hold in the phasor form.

To see this, suppose $v_{1}, v_{2}, \ldots, v_{n}$ are the voltages around a closed loop, then

$$
v_{1}+v_{2}+\cdots+v_{n}=0 .
$$

If each voltage $v_{i}$ is a sinusoid, ie.

$$
v_{i}=V_{m i} \cos \left(\omega t+\phi_{i}\right)=\operatorname{Re}\left\{\mathbf{V}_{i} e^{j \omega t}\right\}
$$

with phasor $\mathbf{V}_{i}=V_{m i} \angle \phi_{i}=V_{m i} e^{j \phi_{i}}$, then

$$
\operatorname{Re}\left\{\left(\mathbf{V}_{1}+\mathbf{V}_{2}+\cdots+\mathbf{V}_{n}\right) e^{j \omega t}\right\}=0
$$

which must be true for all time $t$. To satisfy this, we need
(2) $K v L$

$$
\mathbf{V}_{1}+\mathbf{V}_{2}+\cdots+\mathbf{V}_{n}=0
$$

Hence, KVL holds for phasors.
Similarly, we can show that KCL holds in the frequency domain, i.e., if the currents $i_{1}, i_{2}, \ldots, i_{n}$ are the currents entering or leaving a closed surface at time $t$, then

$$
i_{1}+i_{2}+\cdots+i_{n}=0
$$

If the currents are sinusoids and $\mathbf{I}_{1}, \mathbf{I}_{2}, \ldots, \mathbf{I}_{n}$ are their phasor forms, then

$$
\mathbf{I}_{1}+\mathbf{I}_{2}+\cdots+\mathbf{I}_{n}=0 .
$$

7.5.4. Major Implication: Since Ohm's Law and Kirchoff's Laws hold in phasor domain, all resistance combination formulas, volatge and current divider formulas, analysis methods (nodal and mesh analysis) and circuit theorems (linearity, superposition, source transformation, and Thevenin's and Norton's equivalent circuits) that we have previously studied for dc circuits apply to ac circuits !!!

## Just think of impedance as a complex-valued resistance!!

The three-step analysis in the next chapter is based on this insight.
In addition, our ac circuits can now effortlessly include capacitors and inductors which can be considered as impedances whose values depend on the frequency $\omega$ of the ac sources!!

$$
\vec{V}-\vec{V}_{1}-\vec{V}_{2}-\cdots-\vec{V}_{N}=0
$$

### 7.6. Impedance Combinations

Consider $N$ series-connected impedances as shown below.


The same current I flows through the impedances. Applying KVL around the loop gives

$$
\mathbf{V}=\mathbf{V}_{1}+\mathbf{V}_{2}+\cdots+\mathbf{V}_{N}=\mathbf{I}\left(\mathbf{Z}_{1}+\mathbf{Z}_{2}+\cdots+\mathbf{Z}_{N}\right)
$$

The equivalent impedance at the input terminals is
at terminals

$$
\begin{aligned}
& \text { terminals } \\
& a-b
\end{aligned} \rightarrow \mathbf{Z}_{e q}=\frac{\mathbf{V}}{\mathbf{I}}=\mathbf{Z}_{1}+\mathbf{Z}_{2}+\cdots+\mathbf{Z}_{N}
$$

In particular, if $N=2$, the current through the impedance is


$$
\mathbf{I}=\frac{\mathbf{V}}{\mathbf{Z}_{1}+\mathbf{Z}_{2}}
$$

Because $\mathbf{V}_{1}=\mathbf{Z}_{1} \mathbf{I}$ and $\mathbf{V}_{2}=\mathbf{Z}_{2} \mathbf{I}$,

$$
\mathrm{V}_{1}=\frac{\mathrm{Z}_{1}}{\mathrm{Z}_{1}+\mathrm{Z}_{2}} \mathrm{~V}, \quad \mathrm{~V}_{2}=\frac{\mathrm{Z}_{2}}{\mathrm{Z}_{1}+\mathrm{Z}_{2}} \mathrm{~V}
$$

which is the voltage-division relationship.
In general, $\vec{V}_{k}=\frac{\vec{z}_{n}}{\vec{z}_{1}+\vec{z}_{2}+\cdots+\vec{z}_{N}} \stackrel{\rightharpoonup}{V}$

Now, consider $N$ parallel-connected impedances as shown below.


The voltage across each impedance is the same. Applying KCL at the top node gives

$$
\mathbf{I}=\mathbf{I}_{1}+\mathbf{I}_{2}+\cdots+\mathbf{I}_{N}=\mathbf{V}\left(\frac{1}{\mathbf{Z}_{1}}+\frac{1}{\mathbf{Z}_{2}}+\cdots+\frac{1}{\mathbf{Z}_{N}}\right) .
$$

The equivalent impedance $\mathbf{Z}_{e q}$ can be found from

$$
\frac{1}{\mathbf{Z}_{e q}}=\frac{\mathbf{I}}{\mathbf{V}}=\frac{1}{\mathbf{Z}_{1}}+\frac{1}{\mathbf{Z}_{2}}+\cdots+\frac{1}{\mathbf{Z}_{N}} . \quad \overrightarrow{\mathbf{Z}}_{\text {eq }}=\overline{\frac{1}{\vec{Z}_{1}}+\frac{1}{\vec{Z}_{2}}+\cdots+\frac{1}{\overrightarrow{\mathbf{Z}}_{N}}}
$$

When $N=2$,

$$
\mathbf{Z}_{e q}=\frac{\mathbf{Z}_{1} \mathbf{Z}_{2}}{\mathbf{Z}_{1}+\mathbf{Z}_{2}}
$$

Because

$$
\mathbf{V}=\mathbf{I} \mathbf{Z}_{e q}=\mathbf{I}_{1} \mathbf{Z}_{1}=\mathbf{I}_{2} \mathbf{Z}_{2},
$$

we have

$$
\mathrm{I}_{1}=\frac{\mathrm{Z}_{2}}{\mathrm{Z}_{1}+\mathrm{Z}_{2}} \mathbf{I}, \quad \mathbf{I}_{2}=\frac{\mathrm{Z}_{1}}{\mathrm{Z}_{1}+\mathrm{Z}_{2}} \mathbf{I}
$$

which is the current-division principle.

at terminals $a-b$
Example 7.6.1. Find the input impedance of the circuit below. Assume that the circuit operates at $\omega=50$ rads.


Example 7.6.2. Determine $v_{n}(t)$ in the circuit below.

step 1: conversion to phasor domain



[^0]:    ${ }^{1}$ When there are multiple sources, we assume that all sources are at the same frequency.

